Nonlinear Control of PVTOL Vehicles subjected to Drag and Lift Forces

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Conclusions and future work $_{\rm OOO}$

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Aerial vehicles





Rotary-wings in hovering phase

- Lift forces are negligible
- Thrust counteracts the weight
- High energy consumption

Fixed-wings in cruising phase

- Lift compensates the weight
- Thrust counteracts the drag
- Low energy consumption

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Aerial vehicles





Rotary-wings in hovering phase

• Nonlinear feedback methods developed by neglecting the aerodynamic forces

Fixed-wings in cruising phase

- Feedback methods for linearized aerodynamic models
- Only local stability is concerned

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Convertible vehicles



In vertical take-off phase

- the aerodynamic forces are negligible
- control design similar to the one for hovering phase

Control requirement



In cruising flight phase

- the aerodynamic forces are preponderant
- control design similar to the one for cruising phase

Robust transition maneuvers from take-off to cruising flight

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Objectives

- Analysis and modeling of the aerodynamic reaction forces
- Development of nonlinear controllers by taking into account the nonlinearities of the aerodynamic reaction forces
- Development of a unique robust control law from take-off to cruising flight

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Problem statement, notation and system modeling

A single actuated vehicle immersed in air which exerts reaction forces is considered in the 3-D Lie group SE(2)

• heta angle between $ec{\imath}_0$ and $ec{\imath}$



Assumption

Complete torque actuation $\Longrightarrow \omega$ as control input

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Aerodynamic reaction forces

- Interactions between a solid body and the surrounding fluid are governed by *Navier-Stokes equations*
- They are a set of nonlinear partial differential equations involving
 - viscosity,
 - compressibility,
 - and *density* of the fluid

• Solving the equations requires spatial integration over the shape of the body which typically does not yield closed-form expressions

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Analytical expression of the aerodynamic forces

Lift and drag decomposition

$$\begin{split} F_{a} &= F_{L} + F_{D} \\ F_{L} &= k_{a} |\dot{x}_{a}| c_{L}(\cdot) \dot{x}_{a}^{\perp} \\ F_{D} &= -k_{a} |\dot{x}_{a}| c_{D}(\cdot) \dot{x}_{a} \end{split}$$

$$c_L = c_L(R_e, M, \alpha)$$

 $c_D = c_D(R_e, M, \alpha)$

- $k_a := \frac{1}{2}\rho\Sigma$, ρ is the air density, Σ is the length of the body
- $c_L(\cdot)$ is the *lift coefficient*
- c_D(·) > 0 is the drag coefficient
- R_e is the Reynolds number
- *M* is the *Mach number*
- *α* is the *angle of attack*

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Analytical expression of the aerodynamic forces



Flow regimes: a thumb criterion

- Subsonic flow if M < 0.8.
- *Transonic flow* if 0.8 < *M* < 1.2.
- Supersonic flow if 1.2 < M < 5.
- Hypersonic flow if M > 5.

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Aerodynamic characteristics in the low-subsonic regime

Mach number M < 0.5

- The air ρ is approximately constant
- *M* does not greatly affect c_L and $c_D \Rightarrow$

•
$$c_L = c_L(R_e, \alpha), \ c_D = c_D(R_e, \alpha)$$

• Typical data for c_L and c_D



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Aerodynamic characteristics in the low-subsonic regime





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Single frequency approximation in low-subsonic regime

Assumption: bi-symmetric body

$$\left\{egin{array}{l} c_D(lpha)=c_D(-lpha)\ c_L(lpha)=-c_L(-lpha)\ c_D(lpha)=c_D(lpha+\pi)\ c_L(lpha)=c_L(lpha+\pi)\end{array}
ight.$$

Single frequency approximation

$$\begin{cases} c_D(\alpha) = c_{D_0} + 2c_1 \sin^2(\alpha) \\ c_L(\alpha) = c_1 \sin(2\alpha) \end{cases}$$

Important fact:

the value for c_1 must be properly chosen, i.e. solution to least square minimization problem.

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Single frequency approximation in low-subsonic regime



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Introduction to the control problem

System's dynamics

$$egin{aligned} & m\ddot{x} = - TR(heta)e_1 + mge_1 + F_a(\dot{x}_a, heta) \ & \dot{ heta} &= \omega \end{aligned}$$

Spherical shapes (Hua, Hamel, Morin & Samson 2009)

$$F_a(\dot{x}_a, heta) = F_a(\dot{x}_a) \quad \Rightarrow \quad ext{Control design much simplified}$$

Example. Let $\dot{\tilde{x}} = \dot{x} - \dot{x}_r(t)$. Then $\dot{\tilde{x}} \equiv 0 \iff$

$$-TR(\theta)e_1+F(\dot{x},t)=0$$

 $F(\dot{x}, t)$: "apparent external forces"

The alike-spherical characteristics

Class of functions c_L and c_D for which $T \longrightarrow T_p$ yields

$$m\ddot{x} = -T_{\rho}R(\theta)e_1 + mge_1 + F_{\rho}(\dot{x}_a)$$

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The case of single frequency approximation of c_L and c_D

Assume that the resultant of the aerodynamic forces is of the form $F_a = k_a |\dot{x}_a| \left[c_L(\alpha)S - c_D(\alpha)I \right] \dot{x}_a.$

Main result

$$m\ddot{x} = -TR(\theta)e_1 + mge_1 + F_a(\dot{x}_a, \theta)$$

can be transformed into the form

$$m\ddot{x} = -T_{p}R(\theta)e_{1} + mge_{1} + F_{p}(\dot{x}_{a})$$

with F_p independent of θ , if

$$\left\{egin{array}{l} c_D(lpha)=c_{D_0}+2c_1\sin^2(lpha)\ c_L(lpha)=c_1\sin(2lpha) \end{array}
ight.$$

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Control design

Once the transformation is done, the control design is similar to the one for systems subjected to drag forces only

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A control solution based on (Hua et al. 2009)

- Control laws guarantee (almost)-global stability domain for $x_r(t)$
- Ensures robustness properties



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From hovering to cruising flight



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Conclusions

Absence of the stall phenomena

- Control problem recasted to the one of a spherical shape
- Control design is similar to the one for systems subjected to drag
- Using (Hua et al. 2009) (almost)-global stability and robustness to constant disturbances can be achieved

Stall phenomena

- Existence of equilibrium point is ensured
- Uniqueness of the equilibrium point is not ensured
- Reference trajectories ensuring continuity of $\theta_e(t)$
- Global stability difficult to achieve
- Laws ensuring local stability of x_r or \dot{x}_r have been developed
- Robustness to unmodelled dynamics are achieved

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Future work

- Nature of the equilibrium points
- Positivity of the thrust
- Convergence to the e.p. endowed with the smallest drag
- Enlarge the domain of attraction
- Take into account the ω -dynamics
- 3D-Case ...

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Bibliography			

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